

Thermal Models in High Energy Physics.

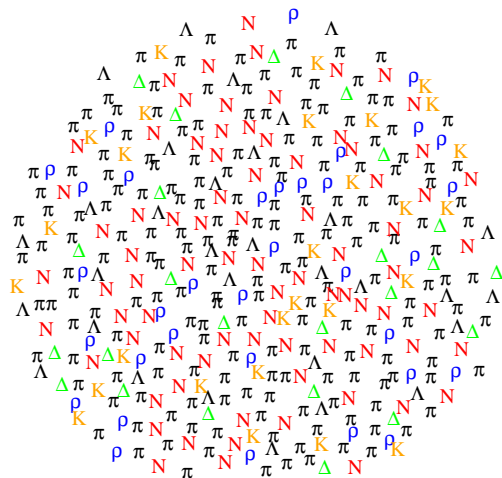
J. Cleymans
University of Cape Town,
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South Africa

**6th EU-Russia-JINR@Dubna Round Table Discussion,
3 - 5 March 2014**

Topics:

- **Thermal Models from 1 to 7000 GeV**
- $\sqrt{s_{NN}} = 11$ GeV is special
- **The Tsallis Distribution at the LHC**

Hadronic Gas



J.C. and H. Satz, Z. fuer Physik C57, 135, 1993.

	Equilibrium
π	$\exp\left[-\frac{E_\pi}{T}\right]$
N	$\exp\left[-\frac{E_N}{T} + \frac{\mu_B}{T}\right]$
\bar{N}	$\exp\left[-\frac{E_N}{T} - \frac{\mu_B}{T}\right]$
Λ	$\exp\left[-\frac{E_\Lambda}{T} + \frac{\mu_B}{T} - \frac{\mu_S}{T}\right]$
$\bar{\Lambda}$	$\exp\left[-\frac{E_\Lambda}{T} - \frac{\mu_B}{T} + \frac{\mu_S}{T}\right]$
K	$\exp\left[-\frac{E_K}{T} + \frac{\mu_S}{T}\right]$
\bar{K}	$\exp\left[-\frac{E_K}{T} - \frac{\mu_S}{T}\right]$

SPS data.

	Measurement
Pb–Pb 158A GeV	
$(\pi^+ + \pi^-)/2.$	600 ± 30
K^+	95 ± 10
K^-	50 ± 5
K_S^0	60 ± 12
p	140 ± 12
\bar{p}	10 ± 1.7
ϕ	7.6 ± 1.1
Ξ^-	4.42 ± 0.31
Ξ^-	0.74 ± 0.04
$\bar{\Lambda}/\Lambda$	0.2 ± 0.04

SPS data.

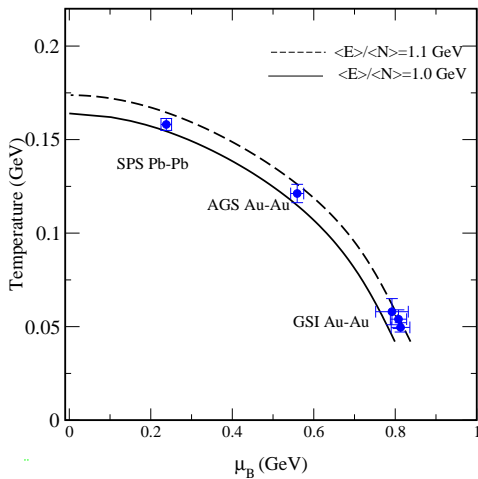
SPS: Freeze-Out Parameters:

$$T = 156.0 \pm 2.4 \text{ MeV}$$

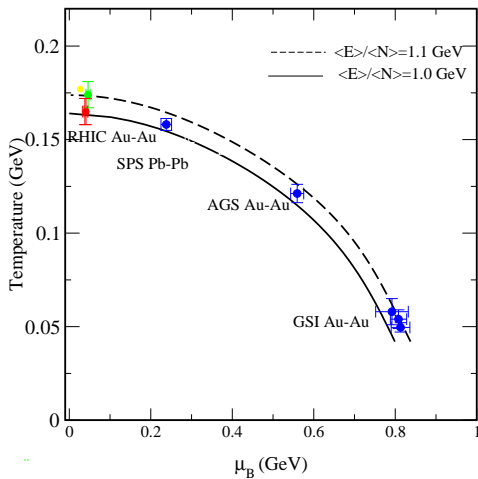
$$\mu_B = 239 \pm 12 \text{ MeV}$$

F. Becattini, J.C., A. Keränen, E. Suhonen and K. Redlich
Physical Review C64 (2001) 024901.

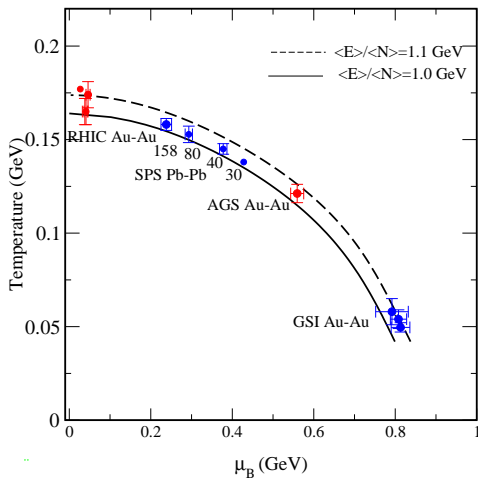
E/N in 1999



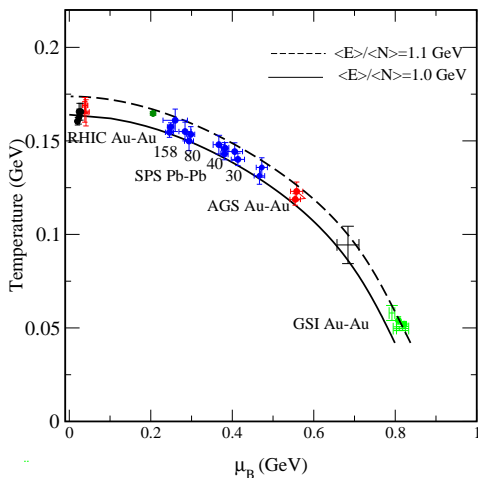
E/N in 2000



E/N in 2005

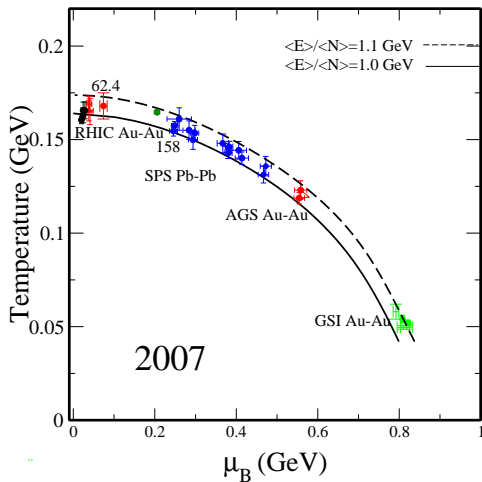


E/N in 2006

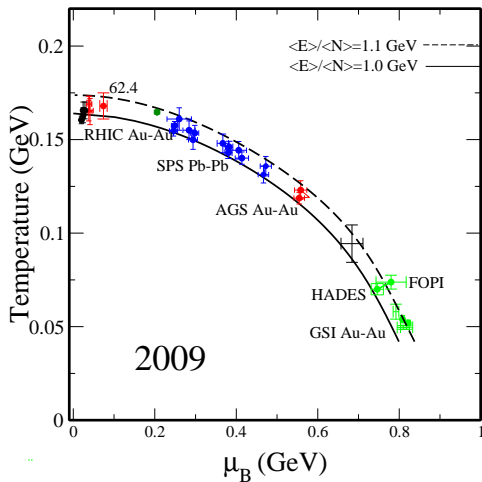


- A. Andronic, P. Braun-Munzinger, J. Stachel, Nucl. Phys. A772, 167, 2006
 J. Manninen, F. Becattini, M. Gazdzicki, Phys. Rev. C73 044905, 2006
 R. Picha, U of Davis, Ph.D. thesis 2002

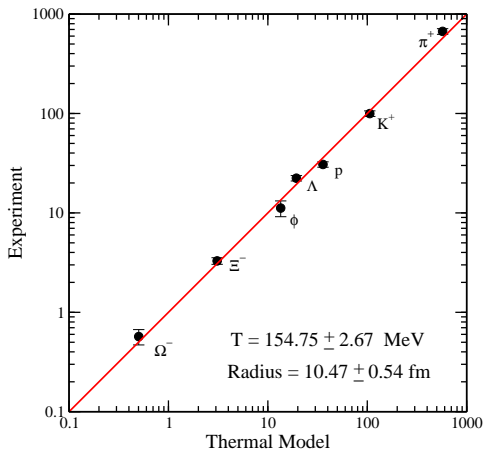
E/N in 2007

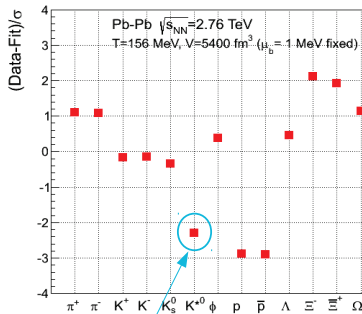
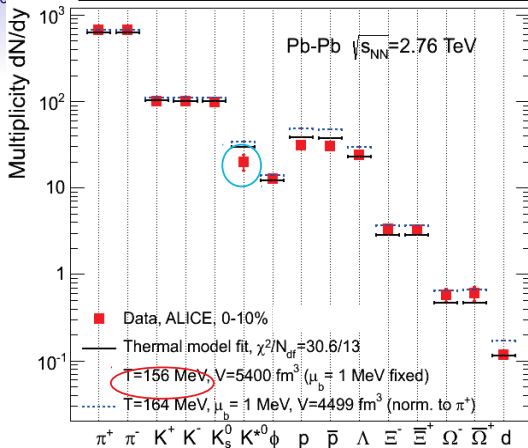


E/N in 2009



2013





strongly decaying
resonance

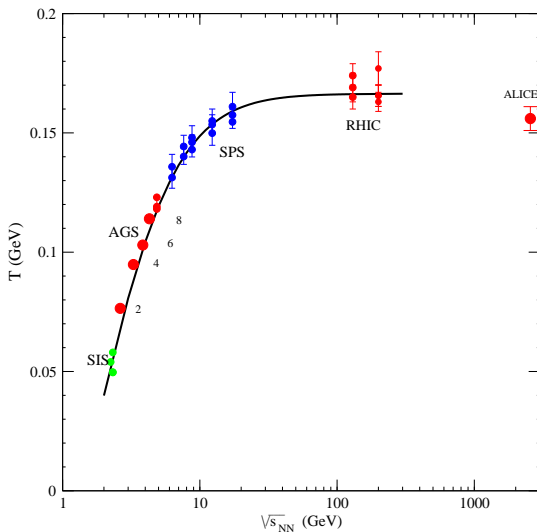
protons low by 2

as compared to 2012: more and final data
T went from 152 to 156 MeV
red. χ^2 went from 4 to 2.35

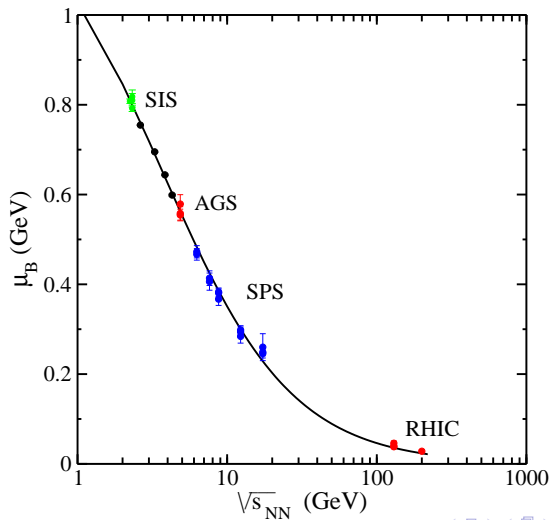
Johanna Stachel

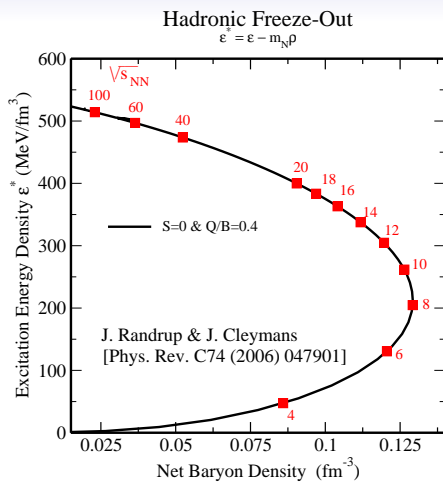
RUPRECHT-KARLS-UNIVERSITÄT HEIDELBERG

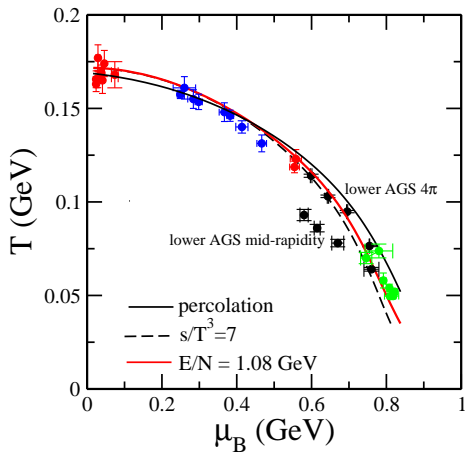
Chemical Freeze-Out Temperature

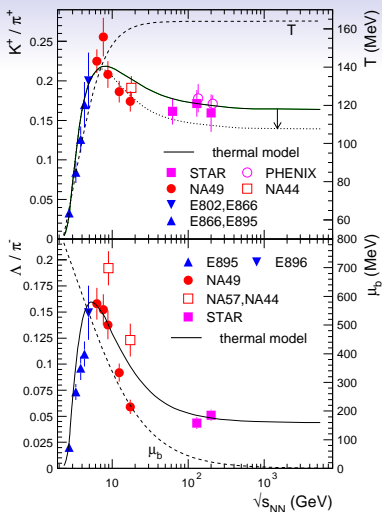


Chemical Freeze-Out μ_B









A. Andronic, P. Braun-Munzinger, J. Stachel, Physics Letters B673 (2009) 142.

In the thermal model the roller-coaster seen in the particle ratios corresponds to a transition from a baryon-dominated to a meson-dominated hadronic gas. This transition occurs at a

- temperature $T = 151$ MeV,
- baryon chemical potential $\mu_B = 327$ MeV,
- energy $\sqrt{s_{NN}} = 11$ GeV.

In the thermal model this transition leads to peaks in the $\Lambda / \langle \pi \rangle$, K^+ / π^+ , Ξ^- / π^+ and Ω^- / π^+ ratios but all occur at slightly different energies.

The maximum in the K^+ / π^+ ratio can be reproduced in thermal models but not the sharpness of the peak.

NEED SMALL STEPS IN BEAM ENERGY!!

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NEED SMALL STEPS IN BEAM ENERGY!!

The Tsallis Distribution at the LHC

Transverse Momentum Distribution

STAR collaboration, B.I. Abelev et al.

arXiv: nucl-ex/0607033; Phys. Rev. **C75**, 064901 (2007)

PHENIX collaboration, A. Adare et al.

arXiv: 1102.0753 [nucl-ex]; Phys. Rev. **C83**, 064903 (2011)

ALICE collaboration, K. Aamodt et al.

arXiv: 1101.4110 [hep-ex]; Eur. Phys. J. **C71**, 1655 (2011)

CMS collaboration, V. Khachatryan et al.

arXiv: 1102.4282 [hep-ex]; JHEP **05**, 064 (2011)

ATLAS collaboration, G. Aad et al.

arXiv: 1012.5104 [hep-ex]; New J. Phys. **13** (2011) 053033.

Transverse Momentum Distribution

STAR, PHENIX, ALICE, CMS, ATLAS use:

$$\frac{d^2N}{dp_T dy} = p_T \times \frac{dN}{dy} \frac{(n-1)(n-2)}{nT(nT + m_0(n-2))} \left(1 + \frac{m_T - m_0}{nT}\right)^{-n}$$

What is the connection with the Tsallis distribution?

Also, the physical significance of the parameters n and T has never been discussed by STAR, PHENIX, ALICE, ATLAS, CMS.

Tsallis Distribution

Possible generalization of Boltzmann-Gibbs statistics

Constantino Tsallis
Rio de Janeiro, TBPf
J. Stat. Phys. 52 (1988) 479-487

Citations: 1 389
Citations in HEP: 513



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MINISTÉRIO DA CIÊNCIA E TECNOLOGIA

**CBPF**

CENTRO BRASILEIRO DE PESQUISAS FÍSICAS

Notas de Física

CBPF-NF-062/87

POSSIBLE GENERALIZATION OF BOLTZMANN-GIBBS
STATISTICS

by

Constantino TSALLIS

RIO DE JANEIRO
1987

Entropy: Tsallis vs Boltzmann

The Boltzmann entropy is given by

$$S^B = -g \sum_i [f_i \ln f_i - f_i], \quad (1)$$

The Tsallis entropy is given by

$$S_T^B = -g \sum_i [f_i^q \ln_q f_i - f_i], \quad (2)$$

which uses

$$\ln_q(x) \equiv \frac{x^{1-q} - 1}{1 - q}, \quad (3)$$

often referred to as q-logarithm.

By maximizing the entropy one obtains expressions for particle density, energy density and pressure.

Maxmizing the Entropy

$$f_j = \left[1 + (q - 1) \frac{E - \mu}{T} \right]^{-\frac{1}{q-1}} \quad (4)$$

For high energy physics a consistent form of Tsallis thermodynamics is:

$$S = -gV \int \frac{d^3p}{(2\pi)^3} [f^q \ln_q f - f], \quad (5)$$

$$N = gV \int \frac{d^3p}{(2\pi)^3} f^q, \quad (6)$$

$$\epsilon = g \int \frac{d^3p}{(2\pi)^3} E f^q, \quad (7)$$

$$P = g \int \frac{d^3p}{(2\pi)^3} \frac{p^2}{3E} f^q. \quad (8)$$

where T and μ are the temperature and the chemical potential, V is the volume and g is the degeneracy factor.

Thermodynamic consistency

$$dE = -pdV + TdS + \mu dN$$

Inserting $E = \epsilon V$, $S = sV$ and $N = nV$ leads to

$$d\epsilon = Tds + \mu dn$$

$$dP = nd\mu + sdT$$

In particular

$$n = \left. \frac{\partial P}{\partial \mu} \right|_T, \quad s = \left. \frac{\partial P}{\partial T} \right|_{\mu}, \quad T = \left. \frac{\partial \epsilon}{\partial s} \right|_n, \quad \mu = \left. \frac{\partial \epsilon}{\partial n} \right|_s.$$

are satisfied.

Transverse Momentum Distributions

In the Tsallis distribution the total number of particles is given by:

$$N = gV \int \frac{d^3p}{(2\pi)^3} \left[1 + (q-1) \frac{E - \mu}{T} \right]^{-\frac{q}{q-1}}.$$

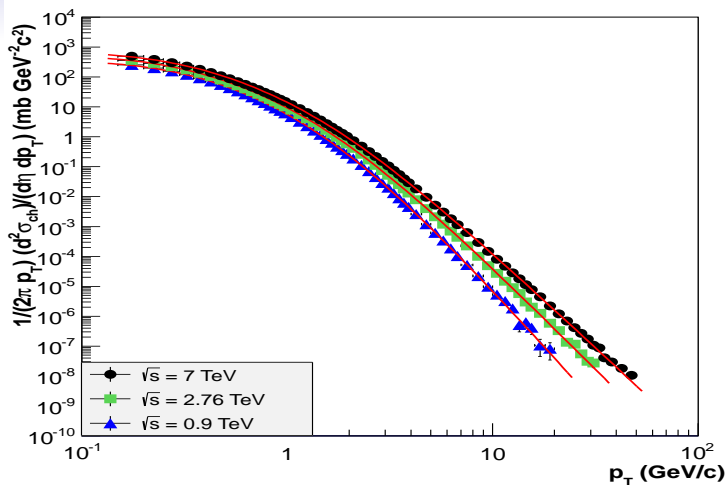
The corresponding momentum distribution is given by

$$E \frac{dN}{d^3p} = gVE \frac{1}{(2\pi)^3} \left[1 + (q-1) \frac{E - \mu}{T} \right]^{-\frac{q}{q-1}},$$

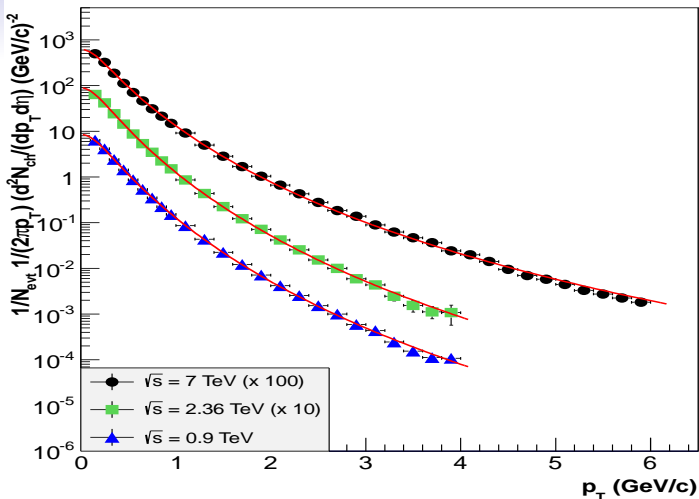
which, in terms of the rapidity and transverse mass variables, $E = m_T \cosh y$, becomes (at mid-rapidity for $\mu = 0$)

$$\left. \frac{d^2N}{dp_T dy} \right|_{y=0} = gV \frac{p_T m_T}{(2\pi)^2} \left[1 + (q-1) \frac{m_T}{T} \right]^{-\frac{q}{q-1}},$$

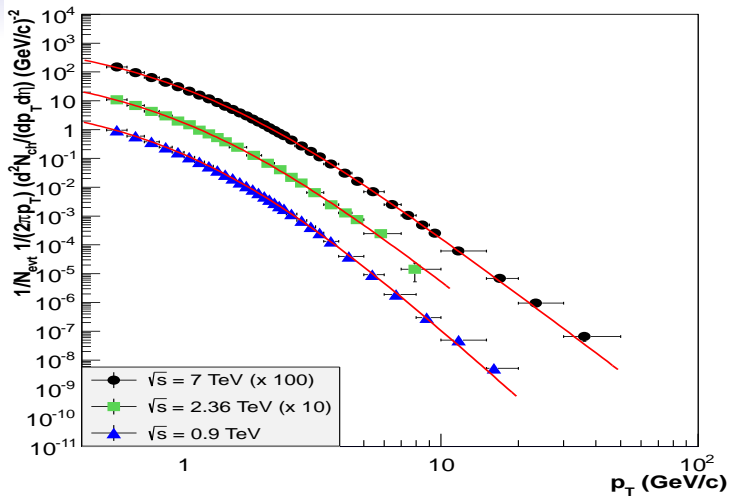
J.C. and D. Worku, J. Phys. **G39** (2012) 025006;
arXiv:1203.4343[hep-ph].



ALICE, Phys. Lett. B **693**, 53 (2010); Eur. Phys. J. C **73**, 2662 (2013).

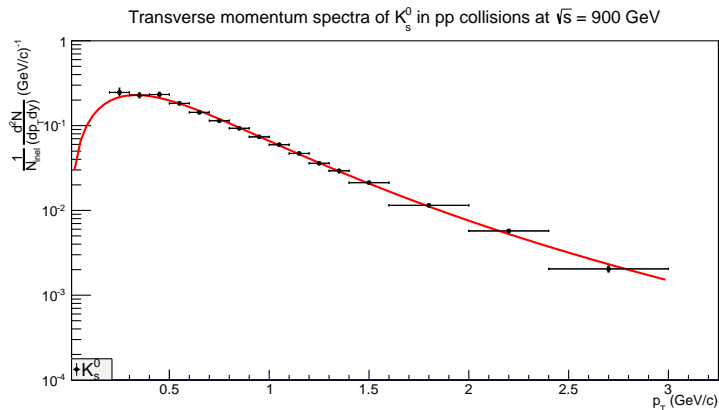


CMS, JHEP **02**, 041 (2010); Phys. Rev. Lett. **105**, 022002 (2010).



ATLAS, New J. Phys. **13**, 053033 (2011).

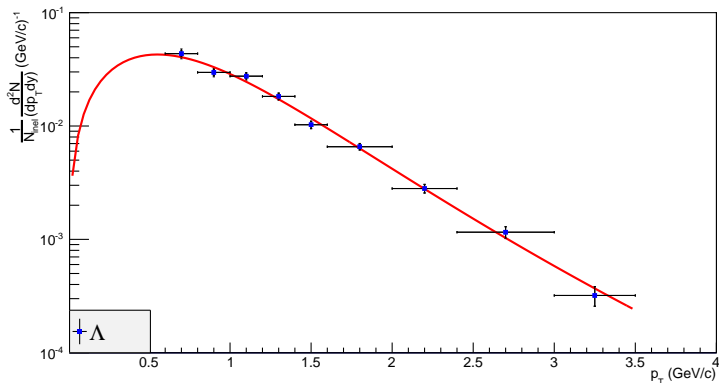
Tsallis Distribution p-p



M.D. Azmi

Tsallis Distribution p-p

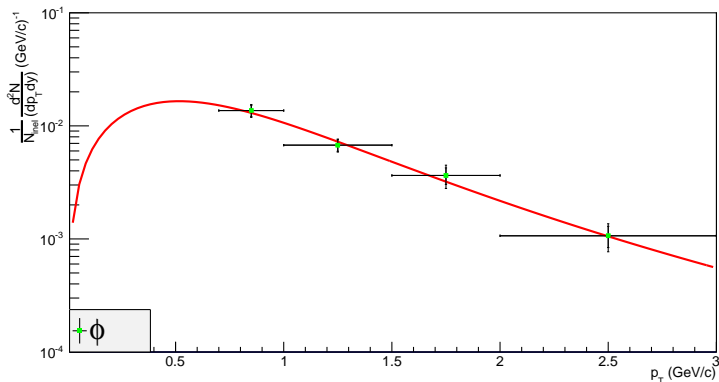
Transverse momentum spectra of Λ in pp collisions at $\sqrt{s} = 900$ GeV



M.D. Azmi

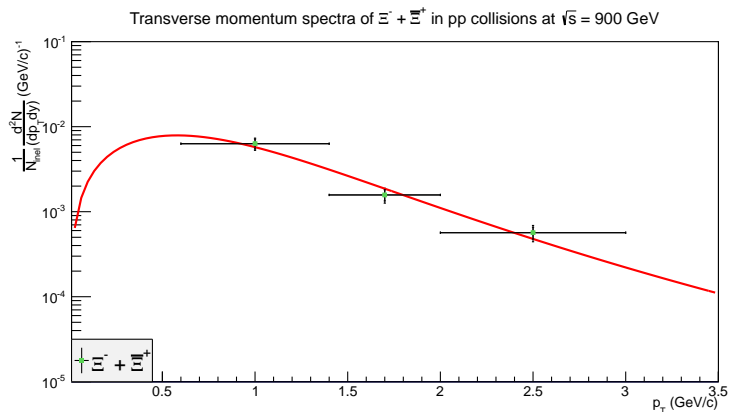
Tsallis Distribution p-p

Transverse momentum spectra of ϕ in pp collisions at $\sqrt{s} = 900$ GeV



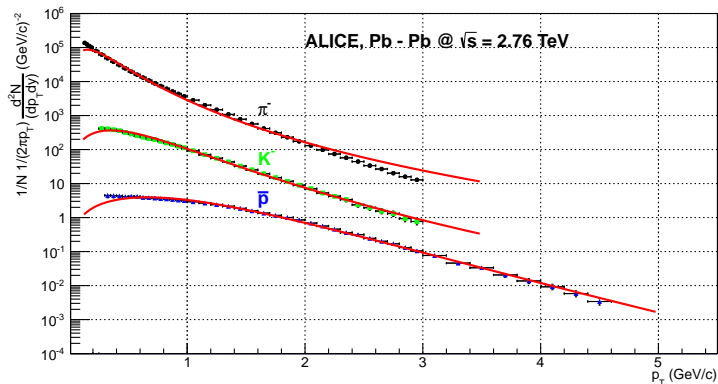
M.D. Azmi

Tsallis Distribution p-p



M.D. Azmi

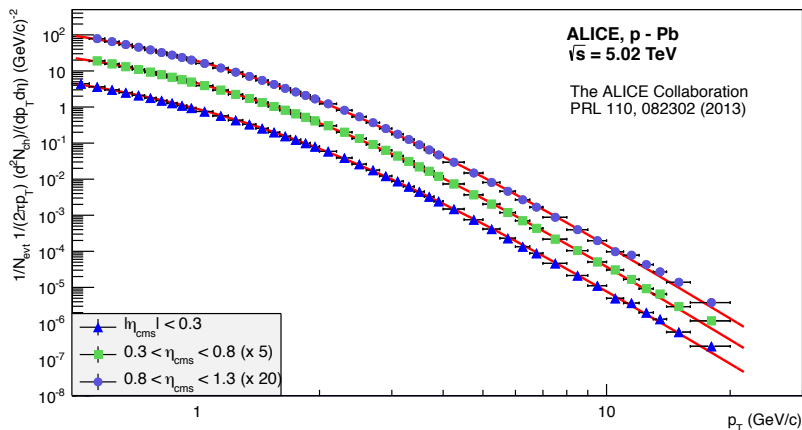
Tsallis Distribution does not describe Pb-Pb



M. Danish Azmi

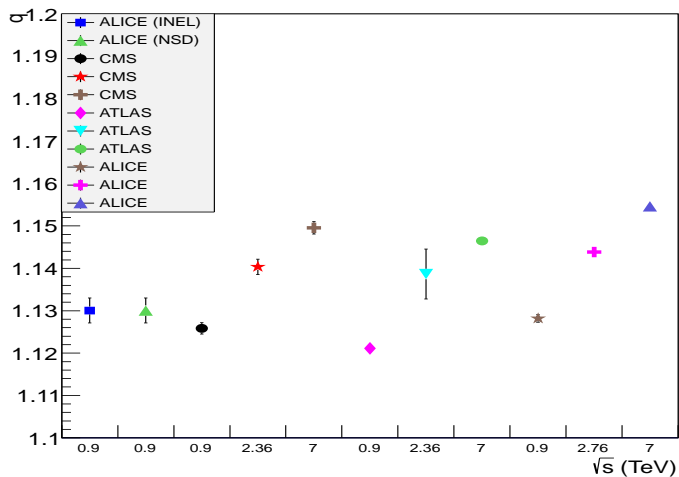
... but works for p-Pb collisions

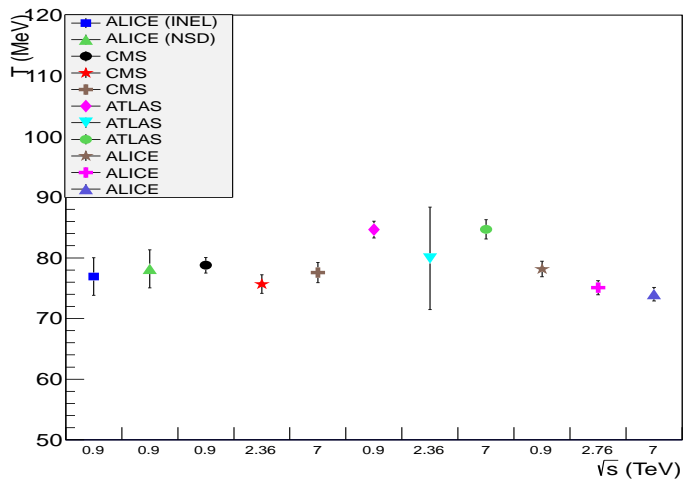
Transverse momentum distribution of charged particles with TSallis Fit in NSD p-Pb collisions

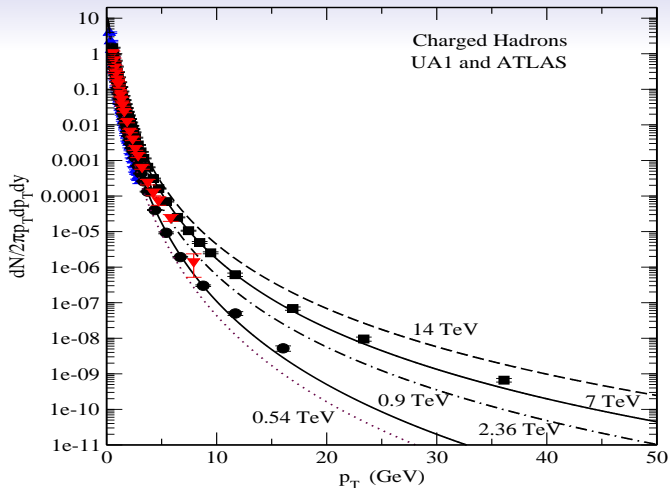


$q = 1.140$

M. Danish Azmi







J.C., G.I. Lykasov, A.S. Sorin, O.V. Teryaev, A.S. Pravan, D. Worku
Physics Letters B 723 (2013) 351-354.

Conclusion:

- Thermal models give a good description of particle yields from 1 to 7000 GeV.
- LHC: Temperature is a bit lower than expected.
- LHC: Too many pions and not enough protons.
- NICA: needed with small steps in beam energy! $\sqrt{s_{NN}} = 11$ GeV is special
- LHC: Final state in Proton-Proton collisions at the LHC produce a system consistent with Tsallis thermodynamics.